

**CONSISTENCY ESTIMATES FOR A DOUBLE-LAYER POTENTIAL
AND APPLICATION TO THE NUMERICAL ANALYSIS
OF THE BOUNDARY-ELEMENT APPROXIMATION
OF ACOUSTIC SCATTERING BY A PENETRABLE OBJECT**

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ABSTRACT. The consistency of the approximation of a double-layer potential when using a boundary-element method and approximating the curved boundary by a polygonal curve in two-dimensional problems or by a polyhedral surface in three-dimensional ones is investigated. The results are applied to the numerical analysis of the approximation of a model problem: the diffraction of a time-harmonic acoustic wave by a penetrable object.

INTRODUCTION

Double-layer potentials are widely used in the numerical computation of solutions to boundary value problems by means of boundary-element methods. Without being exhaustive, we quote: Laplace and Helmholtz equations (cf. Nédélec [24], Giroire [12, 13]), scattering of an electromagnetic wave by a dielectric cylinder (cf. Chang and Harrington [4]), coupling interior and boundary elements (cf. Johnson and Nédélec [18], Sequeira [34]), etc.

In the case of curved boundary, for reasons of effective computation and digital handling of geometrical data, an approximate boundary must be included on which the discrete problem is posed. Although this treatment is quite common in practical computations, as far as we know, no result is available on the consistency or the asymptotic behavior of the error due to this approximation. By contrast, for the single-layer potential, the method initiated by Nédélec [25] gives complete answers in several directions (cf., e.g., Le Roux [22], Giroire [12, 13], Bendali [2, 3], Johnson and Nédélec [18], etc.). This seems to be due to the following feature. The parametrization of the exact boundary is done by an orthogonal projection of the approximate boundary on the former (cf. Nédélec [25]). The consistency error is then essentially obtained by estimating the difference between the potential kernel values for two points on the approximate boundary and for their respective projections on the exact boundary. This could be called "a local technique estimate". Since the singularity of the derivatives of the kernel of single-layer potentials is generally not integrable, this approach is not suited for double-layer potentials. Therefore, a global procedure is needed. The aim of the present work is precisely to deal with this problem. The basic

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idea is to use a Green's formula in the domain bounded by the actual boundary and the approximate one. Then the consistency estimate essentially results in relatively long calculations involving interior elliptic estimates, the so-called inverse property in the finite element method, and exact computation of the potential created by linear distribution on elements of the approximate boundary, initiated by De Hoop (cf. Van Herk [36]).

Among all the boundary value problems whose formulation as boundary integral equation involves a double-layer potential, the results obtained are illustrated by carrying out the numerical analysis of the one which is related to the scattering of a time-harmonic acoustic wave by a penetrable object. Two reasons essentially motivate this choice. The first comes from the fact that the kernel arising in this problem is that of the Helmholtz equation. It is decomposed as a superposition of the kernel of the static problem, which is that of the Laplace equation, and a less singular kernel to which Nédélec's "local technique estimate" applies. Hence, this situation gives answers which can be used in other directions. Beside the fact that this problem is important in itself in the applications, as suggested by the numerous papers devoted to the subject (cf., e.g., Harrington [16], Rokhlin [31], Kleinman and Martin [19], etc.), the second reason is due to the formulation of the problem as a Fredholm system of integral equations of the first kind, free of hypersingular integrals, considered here. Although the formulation as a Fredholm system of integral equations of the second kind is the most commonly used and the easiest to analyze, it is presently well established that the formulation considered in this paper has generally better numerical stability properties (cf. Nédélec [26]).

The formulation as a Fredholm system of the first kind is obtained directly in the two-dimensional case when modeling the scattering of a time-harmonic electromagnetic wave by a dielectric cylinder (cf. Chang and Harrington [4], Clair [6], for example). This formulation appears without the removal of the hypersingular integrals in the paper of Costabel and Stephan [9]. Hypersingular integrals are numerically difficult to handle. The normal trace of a double-layer potential given variationally by Hamdi [15] is the procedure used here to remove the hypersingularities of the integrals from the formulation in both two- and three-dimensional cases. It must be emphasized that the formulation given by Chang and Harrington, starting from Maxwell's equations in the two-dimensional case, directly provides the correct formulation without using Hamdi's formula. It must also be noted that direct numerical approximations of hypersingular integrals recently became available (cf. Schwab and Wendland [32], Krishnasamy et al. [21]).

The paper is organized as follows. In §1 the mathematical statement of the problem of scattering of a time-harmonic acoustic wave by a penetrable obstacle is given. Some known results on existence, uniqueness, and regularity of the solution are recalled. The rest of §1 is devoted to the reduction of the transmission problem to an integro-differential system stated on the boundary of the obstacle. The discretization of this system by a finite element method constitutes the boundary-element method used for the approximation of the solution of the scattering problem. If the consistency error for the approximation of the double-layer potential is not taken into account, the numerical analysis of the problem may be performed by techniques which are by now standard (cf. Nédélec [24], Le Roux [22], etc.). Therefore, §2 of this article is devoted entirely

to this consistency error, and only the final results of the numerical analysis are given.

1. THE SCATTERING PROBLEM AND ITS APPROXIMATION

1.1. The boundary value problem. The acoustic equations (cf., e.g., Germain [11] for general equations and boundary conditions, and Kress and Roach [20] for particular aspects considered here) show that the time-harmonic wave v scattered by a penetrable obstacle is a solution of the following transmission-type boundary value problem.

With the notation

$$(1.1) \quad \begin{cases} \Omega_1 & \text{a bounded } C^\infty \text{ domain of } \mathbb{R}^n; \quad n = 2, 3, \\ \Omega_2 := \mathbb{R}^n \setminus \overline{\Omega}_1 & \text{the exterior domain,} \end{cases}$$

$$(1.2) \quad \begin{cases} \text{find } v \in C^\infty(\overline{\Omega}_1) \cap C^\infty(\overline{\Omega}_2) \text{ such that} \\ \Delta v + k_1^2 v = (k_2^2 - k_1^2)u_0 & \text{in } \Omega_1, \\ \Delta v + k_2^2 v = 0 & \text{in } \Omega_2, \\ [v] = 0 & \text{on } \Gamma, \\ [\chi \partial_{\mathbf{n}} v] = -[\chi \partial_{\mathbf{n}} u_0] & \text{on } \Gamma, \\ \lim_{|x| \rightarrow +\infty} |x|^{(n-1)/2} (\text{grad } v \cdot x/|x| - ik_2 v) = 0. \end{cases}$$

The last relation appearing in (1.2) is the Sommerfeld radiation condition describing the energy propagation from the obstacle towards infinity. In connection with this latter condition, the limit is always assumed to be uniform with respect to the angular variables $x/|x|$.

In the preceding equations (1.2),

- Γ denotes the common boundary to Ω_1 and Ω_2 ,
- \mathbf{n} the unit normal to Γ outwardly directed to Ω_1 ,
- χ a piecewise constant function defined by

$$(1.3) \quad \chi := \begin{cases} 1/\rho_1 & \text{in } \Omega_1, \\ 1/\rho_2 & \text{in } \Omega_2, \end{cases}$$

where $\rho_m > 0$, $m = 1, 2$, is the density of fluid at rest occupying the domain Ω_m ,

u_0 the (given) incident wave (generally a plane wave) satisfying the equation

$$(1.4) \quad \Delta u_0 + k_2^2 u_0 = 0 \quad \text{in a neighborhood of } \overline{\Omega}_1,$$

k_m , $m = 1, 2$, the wave number

$$(1.5) \quad k_m := \omega/c_m,$$

where c_m is the sound velocity in the fluid occupying the domains Ω_m ; the time variation $e^{-i\omega t}$ ($\omega > 0$ is the pulsation) is suppressed by linearity,

$[\cdot]$ is the jump across Γ defined by

$$(1.6) \quad [w] := w^{(1)}|_{\Gamma} - w^{(2)}|_{\Gamma},$$

where $w^{(m)}|_{\Gamma}$, $m = 1, 2$, is the trace on Γ of a function or distribution defined in Ω_m ,

$\mathbf{a} \cdot \mathbf{b} = \sum_{i=1}^n a_i \bar{b}_i$ is the “scalar product” of two vectors respectively with n real or complex components $a_i, b_i, i = 1, \dots, n$.

Other notation, as well as some function spaces like the Schwartz spaces $\mathcal{D}'(\Omega_1), \mathcal{D}(\Omega_1)$, etc., or Sobolev spaces like $H^1(\Omega_1), H^{1/2}(\Gamma)$, etc., are standard in partial differential equation theory (cf., e.g., Lions and Magenes [23], Chazarain and Piriou [5], etc.) and are used without further comment.

It must be emphasized that, as a result of (1.1), Γ is a closed, bounded $(n-1)$ -dimensional C^∞ manifold without boundary, embedded in \mathbb{R}^n . In fact, this amount of smoothness is not really necessary for the subsequent analysis. It seems to be sufficient to assume that Γ is of class C^2 . Since, however, this point is not essential, neither for the main result nor for the techniques which are developed for its proof, we choose not to further complicate the exposition by this discussion.

In order to give an existence-uniqueness theorem for the problem (1.2), and for the purpose of numerical analysis, a problem set in a more general framework is considered.

Given f in $L^2(\Omega_1)$, g in $H^{1/2}(\Gamma)$, and ζ in $H^{-1/2}(\Gamma)$,

$$(1.7) \quad \left\{ \begin{array}{l} \text{find } v \in H^1(\Omega_1) \cap H_{\text{loc}}^1(\bar{\Omega}_2) \text{ such that} \\ \Delta v + k_1^2 v = -f \quad \text{in } \mathcal{D}'(\Omega_1), \\ \Delta v + k_2^2 v = 0 \quad \text{in } \mathcal{D}'(\Omega_2), \\ [v] = g \quad \text{in } H^{1/2}(\Gamma), \\ [\chi \partial_{\mathbf{n}} v] = \zeta \quad \text{in } H^{-1/2}(\Gamma), \\ \lim_{|x| \rightarrow +\infty} |x|^{(n-1)/2} (\text{grad } v \cdot x / |x| - ik_2 v) = 0. \end{array} \right.$$

The behavior of v at infinity is described by the radiation condition and not by its belonging to any Sobolev space (cf. Wilcox [37]). This explains the involvement of the Fréchet space

$$H_{\text{loc}}^1(\bar{\Omega}_2) := \{v \in \mathcal{D}'(\Omega_2) \mid \phi v \in H^1(\Omega_2), \forall \phi \in \mathcal{D}(\mathbb{R}^n)\}$$

in the formulation. For any real number s , the space $H_{\text{loc}}^s(\bar{\Omega}_2)$ is similarly defined.

The following theorem can then be stated.

Theorem 1.1. *The problem (1.7) has one and only one solution. Moreover, for all $s \geq -1$, if the data respectively satisfy $f \in H^{\max(s, 0)}(\Omega_1)$, $g \in H^{s+3/2}(\Gamma)$, and $\zeta \in H^{s+1/2}(\Gamma)$, the following regularity holds:*

$$(1.8) \quad v \in H^{s+2}(\Omega_1) \cap H_{\text{loc}}^{s+2}(\bar{\Omega}_2).$$

Proof. Uniqueness follows by Rellich’s lemma [30] (cf., e.g., Kress and Roach [20]). Existence of a solution can be obtained either by standard arguments involving the limiting absorption principle like in Wilcox [37], or by the elegant method of Lax and Phillips (cf. Phillips [29]). The higher-order regularity, $s > -1$, holds by general results for elliptic systems due to Agmon, Douglis, and Nirenberg [1].

1.2. Integral equations. Let

$$(1.9) \quad u := v + u_0,$$

where v is the scattered wave by the obstacle, solution of problem (1.2), and u_0 is the incident wave. The representation formula for the solution of the Helmholtz equation satisfying an outgoing Sommerfeld radiation condition in the case of an exterior domain (cf., e.g., Colton and Kress [7], Gunter [14]) gives

$$(1.10) \quad u(x) = \begin{cases} \int_{\Gamma} (G_1(x, y)\lambda^{(1)}(y) - \partial_{\mathbf{n}}G_1(x, y)p^{(1)}(y)) d\Gamma(y), & x \in \Omega_1, \\ u_0(x) - \int_{\Gamma} (G_2(x, y)\lambda^{(2)}(y) - \partial_{\mathbf{n}}G_2(x, y)p^{(2)}(y)) d\Gamma(y), & x \in \Omega_2. \end{cases}$$

Here, G_m is the Green's kernel giving the outgoing solutions of the Helmholtz equation,

$$(1.11) \quad G_m(x, y) := \frac{i}{4} \left(\frac{k_m}{2\pi|x-y|} \right)^{(n-2)/2} H_{(n-2)/2}^{(1)}(k_m|x-y|), \quad x \neq y, \quad m = 1, 2,$$

where $H_t^{(1)}$ is the Hankel function of order 1 and index t (cf., e.g., Wilcox [37]), and

$$(1.12) \quad \begin{cases} p^{(m)} := u^{(m)}|_{\Gamma}, \\ \lambda^{(m)} := \partial_{\mathbf{n}}u^{(m)}|_{\Gamma}, \end{cases} \quad m = 1, 2,$$

are the respective first and the second trace of u taken from its values on Ω_m , $m = 1, 2$. Unless mentioned explicitly otherwise, all derivatives occurring under the integral sign are implicitly assumed to be taken with respect to the variable y .

The transmission conditions then show that the determination of u is reduced to that of

$$(1.13) \quad \lambda := \lambda^{(1)}/\rho_1 = \lambda^{(2)}/\rho_2,$$

$$(1.14) \quad p := p^{(1)} = p^{(2)}.$$

The formulae, giving the traces of single- and double-layer potentials, yield

$$(1.15) \quad \begin{cases} \rho_1 V_1 \lambda - N_1 p = p/2, \\ \rho_2 V_2 \lambda - N_2 p = -p/2 + u_0|_{\Gamma}, \\ K_1 \lambda - \frac{1}{\rho_1} D_1 p = \lambda/2, \\ K_2 \lambda - \frac{1}{\rho_2} D_2 p = -\lambda/2 + \frac{1}{\rho_2} \partial_{\mathbf{n}} u_0|_{\Gamma}, \end{cases}$$

where the integral operators involved on Γ are denoted by

$$(1.16) \quad \begin{cases} V_m \lambda(x) := \int_{\Gamma} G_m(x, y)\lambda(y) d\Gamma(y), \\ N_m p(x) := \int_{\Gamma} \partial_{\mathbf{n}_y} G_m(x, y)p(y) d\Gamma(y), \\ K_m \lambda(x) := \int_{\Gamma} \partial_{\mathbf{n}_x} G_m(x, y)\lambda(y) d\Gamma(y), \\ D_m p(x) := \partial_{\mathbf{n}_x} \int_{\Gamma} \partial_{\mathbf{n}_y} G_m(x, y)p(y) d\Gamma(y), \end{cases} \quad \text{for } x \in \Gamma \text{ and } m = 1, 2.$$

It is worth noting that the integrals defining V_m , N_m , and K_m exist as usual Lebesgue integrals, whereas the last operator D_m is hypersingular. Among all the possible methods which can be used to give it a meaning, the easiest and fastest, in the present context, is that in the sense of distributions (cf. Costabel [8] for a more general setting).

The mapping properties of these operators are as follows:

$$(1.17) \quad \begin{cases} V_m: H^{-1/2+s}(\Gamma) \rightarrow H^{1/2+s}(\Gamma), \\ N_m: H^{-1/2+s}(\Gamma) \rightarrow H^{1/2+s}(\Gamma), \\ K_m: H^{-1/2+s}(\Gamma) \rightarrow H^{1/2+s}(\Gamma), \\ D_m: H^{1/2+s}(\Gamma) \rightarrow H^{-1/2+s}(\Gamma) \end{cases}$$

for $m = 1, 2$ and for all real s if Γ is C^∞ (cf., e.g., Chazarain and Piriou [5] and Seeley [33]). Costabel [8] has shown that the first and the last relations of (1.17) remain true if Γ is only Lipschitz for $s < 1/2$.

A Fredholm system of the first kind can be obtained by adding the first to the second equation, and the third to the last one. Then, multiplying by test functions and integrating once more over Γ leads to the following variational system:

$$(1.18) \quad \begin{cases} \text{find } (\lambda, p) \in Z := H^{-1/2}(\Gamma) \times H^{1/2}(\Gamma) \text{ such that} \\ \langle V\lambda, \mu \rangle - \langle Np, \mu \rangle = \langle f, \mu \rangle, \quad \forall \mu \in H^{-1/2}(\Gamma), \\ -\overline{\langle q, K\lambda \rangle} + \overline{\langle q, Dp \rangle} = \overline{\langle q, \zeta \rangle}, \quad \forall q \in H^{1/2}(\Gamma), \end{cases}$$

where

$$(1.19) \quad \begin{cases} \langle \cdot, \cdot \rangle \text{ is the (anti)duality pairing between } H^{1/2}(\Gamma) \text{ and } H^{-1/2}(\Gamma), \\ V := \rho_1 V_1 + \rho_2 V_2, \\ N := N_1 + N_2, \\ K := K_1 + K_2, \\ D := \frac{1}{\rho_1} D_1 + \frac{1}{\rho_2} D_2. \end{cases}$$

This formulation was given by Costabel and Stephan [9]. It needs the computation of the hypersingular integrals which appear when the explicit expression of the operator D is considered. These hypersingular integrals are well understood from a theoretical point of view, thanks to the Calderon-Zygmund theory of singular integrals (cf., e.g., Neri [28]) or to the theory of pseudodifferential operators (cf., e.g., Chazarain and Piriou [5], Seeley [33], etc.). It is only recently, as was pointed out to the authors by the referee, that suitable quadrature schemes for finite part integrals (cf. Schwab and Wendland [32]) or regularization techniques based on Taylor expansion were proposed to treat directly the hypersingular operator D .

Here we choose to use the following recipe introduced by Hamdi [15] (cf. also Nédélec [26]), which essentially is an integration by parts on the boundary and which was used with some success in practical computations (cf., e.g., Chang and Harrington [4], Clair [6], etc.):

$$(1.20) \quad \overline{\langle q, Dp \rangle} = -c(p, q) + d(p, q),$$

$$(1.21) \quad c(p, q) := \int_{\Gamma} \int_{\Gamma} \left(\frac{1}{\rho_1} G_1 + \frac{1}{\rho_2} G_2 \right) (x, y) \operatorname{grad}_{\Gamma} p(y) \times \mathbf{n}(y) \\ \cdot \operatorname{grad}_{\Gamma} \overline{q(x)} \times \mathbf{n}(x) d\Gamma(y) d\Gamma(x),$$

$$(1.22) \quad d(p, q) := \int_{\Gamma} \int_{\Gamma} \left(\frac{k_1^2}{\rho_1} G_1 + \frac{k_2^2}{\rho_2} G_2 \right) (x, y) p(y) \mathbf{n}(y) \\ \cdot \overline{q(x)} \mathbf{n}(x) d\Gamma(y) d\Gamma(x),$$

where the multiplication sign \times denotes the exterior vector product $\mathbf{a} \times \mathbf{b} := a_1 b_2 - a_2 b_1$ in two-dimensional problems and the usual vectorial product in three dimensions, and $\operatorname{grad}_{\Gamma}$ is the tangential gradient of a function defined on Γ . It is worth noting that, in the two-dimensional case, $\operatorname{grad}_{\Gamma} p(y) \times \mathbf{n}(y) = \partial_{\sigma} p(y)$, where σ is the curvilinear abscissa which parametrizes Γ in the counterclockwise direction.

The operator K is the transpose of the operator N ; i.e., it satisfies

$$(1.23) \quad \langle q, K\lambda \rangle = \langle \overline{Nq}, \lambda \rangle, \quad \forall q \in H^{1/2}(\Gamma),$$

where the operator \overline{N} is defined by

$$(1.24) \quad \overline{Nq} = \overline{N\overline{q}}.$$

With the sesquilinear forms

$$(1.25) \quad \begin{cases} a(\lambda, \mu) := \langle V\lambda, \mu \rangle, \\ b(p, \mu) := \langle Np, \mu \rangle, \\ b^*(\lambda, q) := \langle \overline{Nq}, \lambda \rangle, \end{cases}$$

the system (1.18) can be rewritten in the form

$$(1.26) \quad \begin{cases} \text{find } (\lambda, p) \in Z \text{ such that} \\ a(\lambda, \mu) - b(p, \mu) = \langle f, \mu \rangle, \quad \forall \mu \in H^{-1/2}(\Gamma), \\ -b^*(\lambda, q) - c(p, q) + d(p, q) = \langle \overline{q}, \zeta \rangle, \quad \forall q \in H^{1/2}(\Gamma). \end{cases}$$

It is well known by now that the operator associated with the system (1.18) is a strongly elliptic pseudodifferential operator on Γ (cf. Costabel and Wendland [10]). Usual duality properties in Hilbert spaces enable one to identify the space T with the (anti)dual space Z' of Z . Then, the theory of pseudodifferential operators (cf., e.g., Chazarain and Piriou [5], Seeley [33]) can be used to decompose this operator in the form $\Lambda + \Theta$, where Λ is coercive on Z , and to establish that Λ , for all real s , is a continuous mapping

$$(1.27) \quad \Lambda: Z_s \rightarrow T_s$$

and Θ a linear compact map from Z_s into T_s , where $Z_s := H^{-1/2+s}(\Gamma) \times H^{1/2+s}(\Gamma)$ and $T_s := H^{1/2+s}(\Gamma) \times H^{-1/2+s}(\Gamma)$.

The uniqueness of the solution of the system (1.18) (cf., e.g., Costabel and Stephan [9]) and the Fredholm alternative then yield the following theorem.

Theorem 1.2. *The operators Λ and $\Lambda + \Theta$ are (algebraic and topological) isomorphisms from Z_s onto T_s for all real s .*

Remark 1.3. In view of the coerciveness properties given by Nédélec and Planchard [27] in the three-dimensional case, and by Hsiao and Wendland [17] in

the two-dimensional case, it is actually a simple exercise to explicitly write out the system (1.26), up to some compact operators, as two uncoupled coercive variational equations. This leads to a simple and direct treatment of the error coming from the approximation of the boundary (cf. Bendali [3], where this approach is detailed for a similar case).

1.3. The boundary element method. According to Nédélec's ideas [25], the approximation of the curved exact boundary Γ seems to require a complicated approach. It is assumed that there exists a *finite family* \mathcal{F} of local charts (i.e., local coordinate patches) (ω, ψ) , forming an atlas of Γ , and that the boundary Γ can be "triangulated" using this atlas. More precisely, it is supposed that \mathcal{F} is a finite family of triples (ω, D, ψ) , where D is a bounded segment in the two-dimensional case and a polygonal domain in the three-dimensional case having its closure \bar{D} contained in ω , which further fulfills the following conditions:

$$(1.28) \quad \Gamma = \bigcup_{(\omega, D, \psi) \in \mathcal{F}} \psi(\bar{D}),$$

and the $\psi(\bar{D})$ are nonoverlapping and such that no vertex of one $\psi_1(\bar{D}_1)$ lies on the (interior of the) curvilinear edge of another $\psi_2(\bar{D}_2)$.

For all $(\omega, D, \psi) \in \mathcal{F}$, consider a "triangulation" $\mathcal{T}_h(D)$ of the domain D , in the usual meaning of finite element triangulations, with elements being segments in the two-dimensional case and triangles in the three-dimensional case. This leads to a triangulation \mathcal{T}_h^* , a disjoint union of the $\mathcal{T}_h(D)$, which is assumed to satisfy the following uniform regularity condition:

$$(1.29) \quad \begin{cases} h_T := \text{diam}(T) \leq h, \\ h/\rho_T \leq c, \end{cases} \quad \forall T \in \mathcal{T}_h^*;$$

here, ρ_T denotes the diameter of the largest ball which can be inscribed in T . As usual in finite element error estimates, c indicates various constants, not necessarily the same in all instances, but always independent of h and of the solutions of the continuous or the discrete problem.

Let (ω, D, ψ) be in \mathcal{F} and T in $\mathcal{T}_h(D)$, and let ψ_T be the linear interpolant of ψ on T from the values of ψ on the vertices of T ,

$$(1.30) \quad K := \psi_T(T),$$

and Γ_h the approximate boundary defined by

$$(1.31) \quad \Gamma_h := \bigcup_{\mathcal{F}} \bigcup_{T \in \mathcal{T}_h(D)} \psi_T(T).$$

We denote by \mathcal{T}_h the collection of all such elements K .

Following Nédélec [25], assume once again that the different triangulations are compatible in the following sense. With Γ_h a closed polygonal curve in the two-dimensional case, and a polyhedral surface in the three-dimensional case, assume the above elements K satisfy an overlapping condition similar to the one mentioned before. Observe that this construction forces the approximate boundary Γ_h to have its vertices on the exact boundary Γ .

For all K in \mathcal{T}_h , we introduce an orthonormal frame of the linear affine manifold supporting K and let $\xi \in \mathbb{R}^{n-1}$ be the respective coordinates. When restricted to each K in \mathcal{T}_h , the following definitions have an obvious meaning.

A function p defined on K is said to belong to \mathbb{P}_m , the space of all complex-valued polynomials of degree less than or equal to m , if p , as function of the variable ξ , is in \mathbb{P}_m . By $\text{grad}_K p$ we denote the tangent vector field to K obtained in the usual way by differentiating with respect to the variable ξ . In this context, $\text{grad}_{\Gamma_h} p$ will be the vector field tangent to Γ_h , defined almost everywhere by $\text{grad}_{\Gamma_h} p|_K := \text{grad}_K p$ for all K in \mathcal{T}_h^* .

Let

$$(1.32) \quad M_h := \{\lambda \in L^2(\Gamma_h) \mid \lambda_K := \lambda|_K \in \mathbb{P}_0, \forall K \in \mathcal{T}_h\},$$

$$(1.33) \quad V_h := \{q \in C^0(\Gamma_h) \mid q_K := q|_K \in \mathbb{P}_1, \forall K \in \mathcal{T}_h\}.$$

The compatibility conditions made on the triangulation \mathcal{T}_h^* insure that the spaces M_h and V_h have dimensions given by the number of elements of \mathcal{T}_h and the number of vertices of Γ_h , respectively.

We can then consider the approximate sesquilinear forms

$$(1.34) \quad a_h(\lambda_h, \mu_h) := \int_{\Gamma_h} \int_{\Gamma_h} (\rho_1 G_1(x_h, y_h) + \rho_2 G_2(x_h, y_h)) \cdot \lambda_h(y_h) \overline{\mu_h(x_h)} d\Gamma_h(y_h) d\Gamma_h(x_h),$$

for λ_h and μ_h in M_h , and

$$(1.35) \quad b_h(p_h, \mu_h) := \int_{\Gamma_h} \int_{\Gamma_h} (\partial_{\mathbf{n}_h} G_1(x_h, y_h) + \partial_{\mathbf{n}_h} G_2(x_h, y_h)) \cdot p_h(y_h) \overline{\mu_h(x_h)} d\Gamma_h(y_h) d\Gamma_h(x_h),$$

for p_h in V_h and μ_h in M_h , and similar expressions for b_h^* , c_h , and d_h . The approximate (anti)linear forms are introduced as follows:

$$(1.36) \quad \begin{cases} \langle f_h, \mu_h \rangle_h := \int_{\Gamma_h} u_0(x_h) \overline{\mu_h(x_h)} d\Gamma_h(x_h) & \text{for } \mu_h \text{ in } M_h, \\ \langle q_h, \zeta_h \rangle_h := -\frac{1}{\rho_2} \int_{\Gamma_h} \partial_{\mathbf{n}_h} u_0(x_h) \overline{q_h(x_h)} d\Gamma_h(x_h) & \text{for } q_h \text{ in } V_h. \end{cases}$$

In this paper, for the sake of clarity, points lying on the approximate boundary and geometrical quantities related to Γ_h are denoted with the help of an index h to distinguish them from the corresponding objects on the exact boundary Γ . Moreover, there is a precise relation between a point denoted by x on Γ and the point x_h on Γ_h , which is given in what follows.

We endow M_h and V_h with different norms. To this end, we begin by defining transport operators which carry each function defined on Γ_h to a function defined on Γ . Following Nédélec [25], we consider the ε -tubular neighborhood of Γ ,

$$(1.37) \quad \mathcal{V}_\varepsilon := \{x \in \mathbb{R}^n \mid \exists y \in \Gamma \text{ such that } |x - y| < \varepsilon\}.$$

If ε is taken sufficiently small, the mapping $\Gamma \times (-\varepsilon, \varepsilon) \rightarrow \mathcal{V}_\varepsilon$ defined by $(m, t) \mapsto m + t\mathbf{n}(m)$ is a C^∞ diffeomorphism (of manifolds), and the projection $\mathcal{P}: \mathcal{V}_\varepsilon \rightarrow \Gamma$ given by $\mathcal{P}(m + t\mathbf{n}(m)) = m$ an infinitely differentiable map (in the meaning of differentiable maps of manifolds). Hereafter, assuming that such an ε is fixed, there is an h^* such that every h less than h^* gives rise to an approximate boundary Γ_h contained in \mathcal{V}_ε . Then $\mathcal{P}|_{\Gamma_h}$ defines a

homeomorphism from Γ_h onto Γ . In the sequel, it will always be assumed that every approximate boundary considered fulfills the later condition.

Hence, for all K in \mathcal{T}_h ,

$$(1.38) \quad \tilde{K} := \mathcal{P}K.$$

A parametrization of a neighborhood of \tilde{K} can then be obtained by

$$(1.39) \quad x = \mathcal{P}x_h(\xi),$$

where $x_h(\xi)$ is a point (of the linear manifold supporting K) in a neighborhood of K with ξ as coordinate.

The Lebesgue measure on Γ may be written as

$$(1.40) \quad d\Gamma(x) := \sqrt{g_K(\xi)} d\xi, \quad x = \mathcal{P}x_h(\xi).$$

As a result, the transport operators are defined by

$$(1.41) \quad \begin{cases} r_h: M_h \rightarrow L^2(\Gamma), \\ s_h: V_h \rightarrow C^0(\Gamma), \\ (r_h \lambda_h)|_{\tilde{K}}(x) := \frac{1}{g_K(\xi)} \lambda_K, \quad x = \mathcal{P}x_h(\xi), \quad x_h \in K, \quad \forall K \in \mathcal{T}_h, \\ (s_h q_h)|_{\tilde{K}}(x) := q_K(\xi), \quad x = \mathcal{P}x_h(\xi), \quad x_h \in K, \quad \forall K \in \mathcal{T}_h. \end{cases}$$

Remark 1.4. The multiplicative term $1/\sqrt{g_K}$ is introduced to simplify some expressions occurring later. All norms written with or without this factor are uniformly equivalent relative to h .

Clearly, the operators r_h and s_h are one-to-one and can be used to identify respectively V_h with a subspace of $C^0(\Gamma)$ and M_h with a subspace of $L^2(\Gamma)$. Consequently, these spaces can be endowed with various norms

$$(1.42) \quad \|\lambda_h\|_s := \|r_h \lambda_h\|_{s, \Gamma}, \quad \lambda_h \in M_h,$$

for all $s \leq 0$, and

$$(1.43) \quad \|p_h\|_s := \|s_h p_h\|_{s, \Gamma}, \quad p_h \in V_h,$$

for all $s \leq 1$.

Here, $\|\cdot\|_{s, \Gamma}$ is a norm of the Sobolev space $H^s(\Gamma)$ related to the previous atlas $\{(\omega, \psi) | (\omega, D, \psi) \in \mathcal{T}\}$.

All the following estimates are obtained by standard techniques (cf. Nédélec and Planchard [27], Nédélec [24], Le Roux [22], etc.).

Lemma 1.5 (Inverse inequalities). *Under the assumption of uniform regularity of the triangulation \mathcal{T}_h^* , we have*

$$(1.44) \quad \|\lambda_h\|_0 \leq ch^{-1/2} \|\lambda_h\|_{-1/2}, \quad \forall \lambda_h \in M_h,$$

$$(1.45) \quad \|p_h\|_1 \leq ch^{-1/2} \|p_h\|_{1/2}, \quad \forall p_h \in V_h,$$

$$(1.46) \quad \left(\sum_{K \in \mathcal{T}_h} \text{meas}(K) |\text{grad}_K p_K|^2 \right)^{1/2} \leq ch^{-1/2} \|p_h\|_{1/2}, \quad \forall p_h \in V_h.$$

The consistency error of the approximation of the sesquilinear forms $a(\cdot, \cdot)$, $c(\cdot, \cdot)$ and $d(\cdot, \cdot)$ and of that of the (anti)linear forms f and ζ are recalled in the following theorem.

Theorem 1.6. *The following consistency error estimates hold:*

$$(1.47) \quad \begin{cases} |a(r_h \lambda_h, r_h \mu_h) - a_h(\lambda_h, \mu_h)| \leq ch \|\lambda_h\|_{-1/2} \|\mu_h\|_{-1/2}, \\ |c(s_h p_h, s_h q_h) - c_h(p_h, q_h)| \leq ch \|p_h\|_{1/2} \|q_h\|_{1/2}, \\ |d(s_h p_h, s_h q_h) - d_h(p_h, q_h)| \leq ch \|p_h\|_{1/2} \|q_h\|_{1/2}, \\ \langle f, \mu_h \rangle - \langle f_h, \mu_h \rangle_h \leq ch^{3/2} \max_{\alpha=1} \sup_{x \in \mathcal{Z}'_\varepsilon} |\partial^\alpha u_0(x)| \|\mu_h\|_{-1/2}, \\ |\langle s_h q_h, \zeta \rangle - \langle q_h, \zeta_h \rangle_h| \leq ch \max_{|\alpha| \leq 2} \sup_{x \in \mathcal{Z}'_\varepsilon} |\partial^\alpha u_0(x)| \|q_h\|_{1/2}, \end{cases}$$

for all λ_h, μ_h in M_h and all p_h, q_h in V_h .

In the above theorem, the incident wave u_0 is assumed to be defined and regular (at least C^2) in a tubular neighborhood \mathcal{Z}'_ε with $\varepsilon' > \varepsilon$.

Remark 1.7. It must be noticed that Le Roux's and Nédélec's results give the estimates

$$(1.48) \quad |a(r_h \lambda_h, r_h \mu_h) - a_h(\lambda_h, \mu_h)| \leq ch^2 \|\lambda_h\|_0 \|\mu_h\|_0,$$

and only

$$(1.49) \quad |d(s_h p_h, r_h q_h) - d_h(p_h, q_h)| \leq ch \|p_h\|_0 \|q_h\|_0,$$

because of the loss of one convergence order when approximating the normal \mathbf{n} to Γ by the normal \mathbf{n}_h to Γ_h . However, the use of the inverse inequality (1.45), to come back to the norm $\|\cdot\|_{-1/2}$, results in a loss of one convergence order in h and makes the behavior of the consistency error of the two sesquilinear forms similar.

2. CONSISTENCY ERROR ESTIMATE FOR A DOUBLE-LAYER POTENTIAL

This section is almost entirely devoted to establishing the behavior of the consistency error for the approximation of the sesquilinear forms related to the double-layer potential, when the actual boundary is replaced by the approximate one. As it is pointed out in the introduction, this estimate cannot be obtained by the usual "local techniques". In order to make the discussion more straightforward, we decompose the kernel yielding the form $b(\cdot, \cdot)$ as a superposition of a singular part, which can be viewed as the kernel of the static case, and a less singular part, which describes the propagation properties of the solution. Obviously, the estimates related to the transposed form b^* are similar. First, in §2.1, we state our main result and give one of its consequences, the numerical analysis of the scattering problem considered in the previous part. The next subsection is devoted to establishing the consistency estimate for the weakly singular part of the double-layer potential. The proof of the full consistency estimate is done in the last subsection and consists of several steps, each embodied by a lemma.

2.1. The main result and its consequence. The main result of this paper is as follows.

Theorem 2.1. *There exists $h^* > 0$ and a constant c not depending on $0 < h \leq h^*$ such that*

$$(2.1) \quad |b(s_h p_h, r_h \mu_h) - b_h(p_h, \mu_h)| \leq ch^{1/2} \|\mu_h\|_{-1/2} \|p_h\|_{1/2},$$

$$(2.2) \quad |b^*(s_h p_h, r_h \mu_h) - b_h^*(p_h, \mu_h)| \leq ch^{1/2} \|\mu_h\|_{-1/2} \|p_h\|_{1/2},$$

for all p_h in V_h and all μ_h in M_h .

The discrete problem, for which the solution may be effectively computed, is stated as follows:

$$(2.3) \quad \begin{cases} \text{find } (\lambda_h, p_h) \in Z_h := M_h \times V_h \text{ such that} \\ a_h(\lambda_h, q_h) - b_h(p_h, \mu_h) = \langle f_h, \mu_h \rangle_h, \quad \forall \mu_h \in M_h, \\ -b_h^*(\lambda_h, q_h) - c_h(p_h, q_h) + d(p_h, q_h) = \overline{\langle q_h, \zeta_h \rangle}_h, \quad \forall q_h \in V_h, \end{cases}$$

where $\langle f_h, \mu_h \rangle_h$ and $\langle q_h, \zeta_h \rangle_h$ are suitable approximations of the linear maps occurring in the continuous problem defined through the data of f in $H^{1/2}(\Gamma)$ and ζ in $H^{-1/2}(\Gamma)$.

Hence, using general results on the numerical analysis of nonconforming approximations of variational problems by a finite element method (cf., e.g., Strang and Fix [35]), Theorem 2.1 provides the last estimate needed to prove:

Theorem 2.2. *There exists $h^* > 0$ and a constant c not depending on $0 < h < h^*$ such that for all (f_h, ζ_h) in $M'_h \times V'_h$, the discrete problem (2.3) has one and only one solution (λ_h, p_h) that satisfies*

$$(2.4) \quad \begin{aligned} & \|r_h \lambda_h - \lambda\|_{-1/2, \Gamma} + \|s_h p_h - p\|_{1/2, \Gamma} \\ & \leq ch^{1/2} (\|f\|_{1/2, \Gamma} + \|\zeta\|_{-1/2, \Gamma}) + \sup_{\mu_h \in M_h} \frac{|\langle f_h, \mu_h \rangle_h - \langle f, r_h \mu_h \rangle|}{\|\mu_h\|_{-1/2}} \\ & \quad + \sup_{q_h \in V_h} \frac{|\langle q_h, \zeta_h \rangle_h - \langle s_h q_h, \zeta \rangle|}{\|q_h\|_{1/2}}. \end{aligned}$$

In the case of particular f and ζ given through the data of the incident wave u_0 (cf. (1.18)), which is assumed to be at least C^2 in a ε' -tubular neighborhood of the boundary Γ with $\varepsilon' > \varepsilon$, and if the approximate linear forms are given by (1.36), the general result of the previous theorem can be stated in a somewhat more precise form.

Corollary 2.3. *Under the above general assumptions, the following error estimate holds:*

$$(2.5) \quad \|r_h \lambda_h - \lambda\|_{-1/2, \Gamma} + \|s_h p_h - p\|_{1/2, \Gamma} \leq ch^{1/2} \max_{|\alpha| \leq 2} \sup_{x \in \mathcal{E}'_\varepsilon} |\partial^\alpha u_0(x)|.$$

2.2. Consistency error estimate for the weakly singular part. The sesquilinear form $b(\cdot, \cdot)$ is defined by means of the kernel

$$(2.6) \quad G = G_1 + G_2,$$

where G_m , $m = 1, 2$, are given in (1.11). Well-known properties of the Bessel functions show that G can be decomposed in the form

$$(2.7) \quad G = E + \Phi,$$

where the kernel E is given by

$$(2.8) \quad E(x, y) := \frac{2}{\omega_n} F_n(|x - y|), \quad x \neq y,$$

and the function Φ may be written as

$$(2.9) \quad \Phi(x, y) = \begin{cases} u(r) & \text{if the dimension } n = 3, \\ r^2 v(r) \ln r + w(r) & \text{if } n = 2. \end{cases}$$

The constant $\omega_n := 2\pi^{n/2}/\Gamma(n/2)$ is the measure of the unit sphere of \mathbb{R}^n , the function $t \mapsto F_n(t)$ a suitable primitive of $t \mapsto -1/t^{n-1}$, $r := |x - y|$, and u, v , and w are analytic functions.

Let $b^{(0)}$ and $b_h^{(0)}$ be the sesquilinear forms defined respectively on Γ and Γ_h through the respective kernels $\partial_n E(\cdot, \cdot)$ and $\partial_{\mathbf{n}_h} E(\cdot, \cdot)$, and $b^{(1)}$ and $b_h^{(1)}$ the sesquilinear forms obtained in the same way via the kernel Φ .

For the weakly singular part of the kernel, the consistency error estimate is established as follows.

Theorem 2.4. *For all μ_h in M_h and p_h in V_h , the following estimate holds:*

$$(2.10) \quad |b^{(1)}(s_h p_h, r_h \mu_h) - b_h^{(1)}(p_h, \mu_h)| \leq ch^{1/2} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}.$$

Proof. First, we take the difference between the respective values of the kernels of the sesquilinear forms $b^{(1)}$ and $b_h^{(1)}$ and write it as

$$(2.11) \quad \begin{aligned} \partial_n \Phi(x, y) - \partial_{\mathbf{n}_h} \Phi(x_h, y_h) &= g(r) \frac{\mathbf{r}}{r} \cdot \mathbf{n}(y) - g(r_h) \frac{\mathbf{r}_h}{r_h} \cdot \mathbf{n}_h(y_h) \\ &= (g(r) - g(r_h)) \frac{\mathbf{r}}{r} \cdot \mathbf{n}(y) + g(r_h) \frac{\mathbf{r}}{r} \cdot (\mathbf{n}(y) - \mathbf{n}_h(y_h)) \\ &\quad + g(r_h) \left(\frac{\mathbf{r}}{r} - \frac{\mathbf{r}_h}{r_h} \right) \cdot \mathbf{n}_h(y_h), \end{aligned}$$

where $\mathbf{r} := x - y$, $r := |\mathbf{r}|$, $\mathbf{r}_h := x_h - y_h$, $r_h := |\mathbf{r}_h|$, and g represents a Lipschitz function over each bounded set of \mathbb{R} .

Hereafter, whenever there are two variables designated by the same letter, the second being indexed by h , such as x and x_h , it is assumed without further comment that the two points x and x_h lie respectively on Γ and Γ_h and that they are related to each other by $x = \mathcal{P}x_h$.

Since x and y are varying on Γ , one easily obtains (cf. Nédélec [25])

$$(2.12) \quad ||x - y| - |x_h - y_h|| \leq ch^2 |x - y| \quad \text{for all } x, y \text{ on } \Gamma.$$

It follows that

$$(2.13) \quad |(g(r) - g(r_h)) \frac{\mathbf{r}}{r} \cdot \mathbf{n}(y)| \leq ch^2 |x - y|.$$

Therefore, the respective part of the expression for $b^{(1)}(s_h p_h, r_h \mu_h) - b_h^{(1)}(p_h, \mu_h)$ can be bounded by $ch^2 \|p_h\|_0 \|\mu_h\|_0$.

Using the estimates

$$(2.14) \quad \begin{cases} |\mathbf{n}(y) - \mathbf{n}_h(y_h)| \leq ch, \\ \left| \frac{\mathbf{r}}{r} - \frac{\mathbf{r}_h}{r_h} \right| \leq ch, \end{cases}$$

one can bound the remaining expressions in (2.11) to obtain terms in $ch \|p_h\|_0 \|\mu_h\|_0$. The inverse inequality (1.44) leads to the estimate (2.10). \square

2.3. Consistency error estimate for the singular part of the kernel. Clearly, the proof of Theorem 2.1 is now reduced to that of

Theorem 2.5. *For all μ_h in M_h and p_h in V_h , the following estimate holds:*

$$(2.15) \quad |b^{(0)}(s_h p_h, r_h \mu_h) - b_h^{(0)}(p_h, \mu_h)| \leq ch^{1/2} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}.$$

Proof. The error to be estimated is denoted by

$$(2.16) \quad e_h := b^{(0)}(s_h p_h, r_h \mu_h) - b_h^{(0)}(p_h, \mu_h), \quad \mu_h \in M_h, p_h \in V_h.$$

Recall that the kernel of $b^{(0)}$ (resp. $b_h^{(0)}$) is $\partial_{\mathbf{n}} E(x, y)$ (resp. $\partial_{\mathbf{n}_h} E(x_h, y_h)$).

It is easy to see that the earlier procedure used for the weakly singular part of the kernel fails here because of the nonconvergence of certain integrals that arise.

For the sake of conciseness, throughout this section the following notation is adopted:

$$(2.17) \quad \begin{cases} p := s_h p_h, & p_h \in V_h, \\ \mu := r_h \mu_h, & \mu_h \in M_h. \end{cases}$$

The error e_h can be written as

$$(2.18) \quad \begin{aligned} e_h = & \int_{\Gamma} \int_{\Gamma} \partial_{\mathbf{n}} E(x, y) p(y) \overline{\mu(x)} d\Gamma(y) d\Gamma(x) \\ & - \int_{\Gamma_h} \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x_h, y_h) p_h(y_h) \overline{\mu_h(x_h)} d\Gamma_h(y_h) d\Gamma_h(x_h). \end{aligned}$$

Recall that h is supposed to be sufficiently small so that Γ_h is contained in the ε -tubular neighborhood \mathcal{V}_ε . All p_h in V_h can be extended to \mathcal{V}_ε in the following way:

$$(2.19) \quad p(\tilde{y}) = p_h(y_h)$$

for all $\tilde{y} = y + t\mathbf{n}(y)$ in \mathcal{V}_ε with $y = \mathcal{P}\tilde{y} = \mathcal{P}y_h$.

This extension gives a function p constant along the normal \mathbf{n} to Γ and coinciding with $s_h p_h$ on Γ and with p_h on Γ_h . Note that this notation is compatible with the one given in (2.17).

In order to avoid unnecessary complications in the proofs and long discussions of every situation that may occur, in the sequel the interior domain Ω_1 is assumed to be strictly convex. However, in other situations the proofs can be adapted to yield the same conclusions.

The remaining part of the proof will be given in a sequence of lemmas.

Lemma 2.6. *Using previously introduced notations in (2.17) and (2.8), the error e_h may be written as*

$$(2.20) \quad \begin{aligned} e_h = & \frac{2}{\omega_n} \int_{\Gamma} (F_n(|\cdot|) * T)(x) \overline{\mu(x)} d\Gamma(x) \\ & + \sum_{K \in \mathcal{T}_h} \int_K \left(\int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x, y_h) - \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x_h, y_h) \right) \\ & \cdot (p(y_h) - p(x_h)) d\Gamma_h(y_h) \overline{\mu_K} dK(x_h), \end{aligned}$$

where T is the compactly supported distribution of \mathbb{R}^n given by

$$(2.21) \quad \langle T, \Phi \rangle = \int_{\Omega_1 \setminus \tilde{\Omega}_1^i} \text{grad } p \cdot \text{grad } \phi \, dx, \quad \forall \phi \in \mathcal{D}(\mathbb{R}^n)$$

and Ω_1^h the bounded domain enclosed by Γ_h .

Proof. Using the definition of $\mu|_{\tilde{K}}$, we can write the error e_h as

$$(2.22) \quad e_h = \sum_{K \in \mathcal{T}_h} \int_K \left(\int_{\Gamma} \partial_{\mathbf{n}} E(x, y) p(y) \, d\Gamma(y) - \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x_h, y_h) p(y_h) \, d\Gamma_h(y_h) \right) \bar{\mu}_K \, dK(x_h).$$

For almost all x in Γ (i.e., such that x_h is in the interior of any K in \mathcal{T}_h), adding and subtracting in (2.22) the term $\int_{\Gamma_k} \partial_{\mathbf{n}_k} E(x, y_h) p(y_h) \, d\Gamma_h(y_h)$ leads to the following expression of e_h :

$$(2.23) \quad e_h = \sum_{K \in \mathcal{T}_h} \int_K \left\{ \int_{\Gamma} \partial_{\mathbf{n}} E(x, y) p(y) \, d\Gamma(y) - \int_{\Gamma_h} \partial_{\mathbf{n}_k} E(x, y_h) p(y_h) \, d\Gamma_h(y_h) \right\} \bar{\mu}_K \, dK(x_h) + \sum_{K \in \mathcal{T}_h} \int_K \left\{ \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x, y_h) p(y_h) \, d\Gamma_h(y_h) - \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x_h, y_h) p(y_h) \, d\Gamma_h(y_h) \right\} \bar{\mu}_K \, dK(x_h).$$

In order to transform the term

$$(2.24) \quad \int_{\Gamma} \partial_{\mathbf{n}} E(x, y) p(y) \, d\Gamma(y) - \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x, y_h) p(y_h) \, d\Gamma_h(y_h) = \sum_{L \in \mathcal{T}_h} \left\{ \int_{\tilde{L}} \partial_{\mathbf{n}} E(x, y) p(y) \, d\tilde{L}(y) - \int_L \partial_{\mathbf{n}_h} E(x, y_h) p(y_h) \, dL(y_h) \right\},$$

the essential tool is to use a Green's formula applied to the functions $E(x, \cdot)$ and p in the open set bounded by the exact and the approximate boundaries Γ and Γ_h . In the previous formula (2.24), $\tilde{L} = \mathcal{P}L$ is a curvilinear element on Γ associated with L in the same way as \tilde{K} was associated with K in (1.38). However, such a treatment cannot be carried out directly because of the lack of regularity of the function p on the latter domain in the three-dimensional case.

Consequently, the Green's formula is applied to the respective functions in each domain U_L bounded by L, \tilde{L} and, in the three-dimensional case, in the faces F_L generated when the normal \mathbf{n} to Γ runs over the boundary $\partial\tilde{L}$ of \tilde{L} ; see Figure 1.

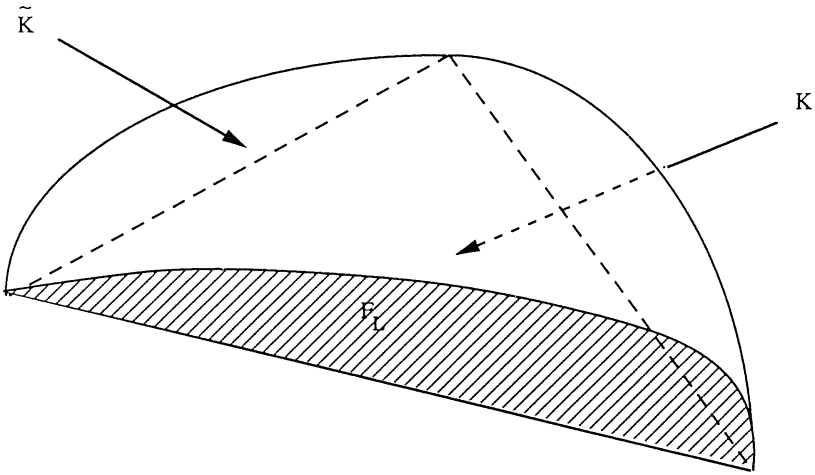


FIGURE 1

Two cases must be distinguished.

(i) The element L is different from the element K . In this case, x does not lie on \tilde{L} . One gets

$$\begin{aligned}
 (2.25) \quad & \int_{U_L} (\Delta E(x, y)p(y) - E(x, y)\Delta p(y)) dy \\
 &= \int_{\tilde{L}} (\partial_{\mathbf{n}} E(x, y)p(y) - E(x, y)\partial_{\mathbf{n}} p(y)) d\tilde{L}(y) \\
 &\quad - \int_L (\partial_{\mathbf{n}_h} E(x, y_h)p(y_h) - E(x, y_h)\partial_{\mathbf{n}_h} p(y_h)) dL(y_h) \\
 &\quad + \int_{F_L} (\partial_{\nu} E(x, y)p(y) - E(x, y)\partial_{\nu} p(y)) dF_L(y),
 \end{aligned}$$

where ν is the unit normal to F_L outwardly directed to U_L . Since p is constant along the normal \mathbf{n} , the equation (2.25) can be reduced to

$$\begin{aligned}
 (2.26) \quad & - \int_{U_L} E(x, y)\Delta p(y) dy = \int_{\tilde{L}} \partial_{\mathbf{n}} E(x, y)p(y) d\tilde{L}(y) + A_L(x) \\
 &\quad - \int_L \partial_{\mathbf{n}_h} E(x, y_h)p(y_h) dL(y_h) \\
 &\quad + \int_{F_L} \partial_{\nu} E(x, y)p(y) dF_L(y),
 \end{aligned}$$

where $A_L(x)$ is defined by

$$(2.27) \quad A_L(x) := \int_L E(x, y_h)\partial_{\mathbf{n}_h} p(y_h) dL(y_h) - \int_{F_L} E(x, y)\partial_{\nu} p(y) dF_L(y).$$

(ii) The element L and the element K are the same element. In this case, the singularity of the function $E(x, \cdot)$ at the point x forces the removal from U_L of all points y such that $|x - y| < \varepsilon$. Taking the limit as ε goes to zero

only has the effect of adding to the right-hand side of (2.26) the term

$$(2.28) \quad \begin{aligned} & \lim_{\varepsilon \rightarrow 0} \int_{S_\varepsilon} \operatorname{grad} E(x, y) \cdot \frac{x-y}{|x-y|} p(y) dS_\varepsilon(y) \\ &= - \lim_{\varepsilon \rightarrow 0} \frac{\omega_n}{2} \frac{2}{\omega_n} F'_n \varepsilon^{n-1} p(x_\varepsilon) = p(x), \end{aligned}$$

where x_ε is a point in $S_\varepsilon := \{y \in \overline{U_L} \mid |x-y| = \varepsilon\}$ whose existence follows from the mean value theorem.

Let P be the function in $L^2(\mathbb{R}^n)$ defined almost everywhere by

$$(2.29) \quad P := \begin{cases} \Delta p(y), & y \in U_L \text{ and } L \in \mathcal{F}_h, \\ 0, & \text{otherwise.} \end{cases}$$

Summing over all L in \mathcal{F}_h , we obtain

$$(2.30) \quad \begin{aligned} - \int_{\mathbb{R}^n} E(x, y) P(y) dy &= \int_{\Gamma} \partial_{\mathbf{n}} E(x, y) p(y) d\Gamma(y) + \sum_{L \in \mathcal{F}_h} A_L(x) \\ &\quad - \int_{\Gamma_h} \partial_{\mathbf{n}_h} E(x, y_h) p(y_h) d\Gamma_h(y_h) + p(x), \end{aligned}$$

since

$$\sum_{L \in \mathcal{F}_h} \int_{F_L} (\partial_{\mathbf{v}} E(x, y) p(y)) dF_L(y) = 0,$$

because p is continuous across F_L and $E(x, \cdot)$ is C^∞ except at the point x .

Thus, using (2.23) and (2.30), we obtain

$$(2.31) \quad \begin{aligned} e_h &= - \int_{\Gamma} \overline{\mu(x)} \sum_{L \in \mathcal{F}_h} A_L(x) d\Gamma(x) - \int_{\Gamma} \overline{\mu(x)} p(x) d\Gamma(x) \\ &\quad + \sum_{K \in \mathcal{F}_h} \int_K \int_{\Gamma_h} (\partial_{\mathbf{n}_h} E(x, y_h) p(y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h)) p(y_h) d\Gamma_h(y_h) \overline{\mu_K} dK(x_h) \\ &\quad - \int_{\Gamma} \overline{\mu(x)} \left(\int_{\mathbb{R}^n} E(x, y) P(y) dy \right) d\Gamma(x). \end{aligned}$$

Using Gauss's theorem, we can write

$$(2.32) \quad p(x_h) = \int_{\Gamma_h} (\partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h)) p(x_h) d\Gamma_h(y_h).$$

This completes the proof of the lemma, since formula (2.20) is now a straightforward consequence of the theory of distributions. \square

In order to bound e_h , the following lemma is proved.

Lemma 2.7. *The distribution T defined in (2.21) satisfies the estimate*

$$(2.33) \quad |(T, \phi)| \leq ch |\operatorname{grad} p|_{0, \Gamma} \|\phi\|_{1, \mathbb{R}^n}, \quad \forall \phi \in \mathcal{D}(\mathbb{R}^n).$$

Proof. Let L be in \mathcal{F}_h . Using the local orthonormal frame associated with L and the parametrization of \tilde{L} given by

$$(2.34) \quad \xi \mapsto \mathcal{P}x_h(\xi),$$

we can express U_L as

$$(2.35) \quad U_L = \{(\xi, t) \in \mathbb{R}^n \mid x_h(\xi) \in L \text{ and } f_L(\xi) < t < 0\}.$$

Here, f_L is a C^∞ function vanishing at all ξ for which $x_h(\xi)$ is a vertex of L .

Using the coordinate system given in an open set of the ε -tubular neighborhood \mathcal{Z}_ε of Γ by

$$(2.36) \quad x(\xi, t) = y(\xi) + t\mathbf{n}(\xi),$$

we can write

$$(2.37) \quad \int_{U_L} \text{grad } p \cdot \text{grad } \phi \, dU_L = \int_L \int_{f_L(\xi)}^0 \sqrt{g_L(\xi, t)} g_L^{\alpha\beta}(\xi, t) \partial_{\xi^\alpha} p \partial_{\xi^\beta} \phi \, d\xi \, dt.$$

Greek indices run from 1 to $n-1$. The contravariant components $g_L^{\alpha\beta}(\xi, t)$ of the metric tensor in the above integral are those related to the coordinate system (ξ, t) for which the Lebesgue measure dU_L is expressed by $dU_L = \sqrt{g_L(\xi, t)} \, d\xi \, dt$. Observe that in (2.37) we use the fact that

$$(2.38) \quad g_L^{\alpha n}(\xi, t) = g_{\alpha n}^L(\xi, t) = 0, \quad \partial_t p = 0.$$

We adopt here the usual convention of tensorial calculus to sum over repeated upper and lower indices.

Usual estimates in the finite element method give $|f_L(\xi)| \leq ch^2$. Since p is not depending on t , one has

$$(2.39) \quad \left(\int_L \int_{f_L(\xi)}^0 |\partial_\gamma p|^2 \, d\xi \, dt \right)^{1/2} \leq ch \left(\int_L |\partial_\gamma p|^2 \, d\xi \right)^{1/2}.$$

The discrete Cauchy-Schwarz inequality completes the proof. \square

Let us decompose the error e_h as $e_h = e_h^1 + e_h^2$, where e_h^1 is the first term in the right-hand side of (2.20). For e_h^1 , we have

Lemma 2.8. *In the above notation,*

$$(2.40) \quad \frac{2}{\omega_n} \left| \int_\Gamma (F_n(|\cdot|) * T)(x) \bar{\mu}(x) \, d\Gamma(x) \right| \leq ch^{1/2} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}$$

for all μ_h in M_h and p_h in V_h .

Proof. Since T is a compactly supported distribution in $H^{-1}(\mathbb{R}^n)$ and the convolution operator $F_n(|\cdot|) *$ is a pseudodifferential operator of order -2 , the following result holds:

$$(2.41) \quad F_n(|\cdot|) * T \in H_{\text{loc}}^1(\mathbb{R}^n).$$

Trace theorems in Sobolev spaces give

$$(2.42) \quad \left| \int_\Gamma \bar{\mu}(x) (F_n(|\cdot|) * T)(x) \, d\Gamma(x) \right| \leq c \|F_n(|\cdot|) * T\|_{1/2, \Gamma} \|\mu_h\|_{-1/2} \\ \leq ch \|\text{grad } p\|_{0, \Gamma} \|\mu_h\|_{-1/2}.$$

By the inverse inequalities of Lemma 1.5, one gets (2.40), proving the lemma. \square

It remains to estimate the term e_h^2 . We decompose it into three terms,

$$(2.43) \quad e_h^2 = e_h^{2,1} + e_h^{2,2} + e_h^{2,3},$$

where we have set

$$(2.44) \quad e_h^{2,1} := \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{N}_K} \int_K \int_L (\partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h)) \cdot (p(y_h) - p(x_h)) \bar{\mu}_K dL(y_h) dK(x_h),$$

$$(2.45) \quad e_h^{2,2} := \sum_{K \in \mathcal{T}_h} \int_K \int_K (\partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h)) \cdot (p(y_h) - p(x_h)) \bar{\mu}_K dK(y_h) dK(x_h),$$

$$(2.46) \quad e_h^{2,3} := \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{A}_K} \int_K \int_L (\partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h)) \cdot (p(y_h) - p(x_h)) \bar{\mu}_K dL(y_h) dK(x_h),$$

with the notations

$$(2.47) \quad \mathcal{N}_K := \{L \in \mathcal{T}_h \mid L \cap K = \emptyset\},$$

$$(2.48) \quad \mathcal{A}_K := \{L \in \mathcal{T}_h \mid L \cap K \neq \emptyset \text{ and } L \neq K\}.$$

Note that, thanks to the uniform regularity of the triangulation, as assumed in (1.29), there exists a constant $c > 0$ not depending on h , L and K such that

$$(2.49) \quad \inf\{|x - y|; x \in K \text{ and } y \in L\} \geq ch, \quad \forall L \in \mathcal{N}_K.$$

The estimation of the term $e_h^{2,1}$ requires the following lemma.

Lemma 2.9. *For all K in \mathcal{T}_h and all L in \mathcal{N}_K , the following estimate holds:*

$$(2.50) \quad |(\partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h))| \leq \frac{ch}{|x - y_h|^{n-1}}$$

for all (x_h, y_h) in $K \times L$.

Proof. Set

$$(2.51) \quad \begin{cases} \mathbf{R} := y_h - x, & R := |\mathbf{R}|, \\ \mathbf{R}_h := y_h - x_h, & R_h := |\mathbf{R}_h|. \end{cases}$$

There follows

$$(2.52) \quad \partial_{\mathbf{n}_h} E(x, y_h) - \partial_{\mathbf{n}_h} E(x_h, y_h) = -\frac{2}{\omega_n} \left(\frac{\mathbf{R}}{R^n} - \frac{\mathbf{R}_h}{R_h^n} \right) \cdot \mathbf{n}_L,$$

where \mathbf{n}_L is the unit normal to the element L coinciding with \mathbf{n}_h on L .

Note that one can write

$$(2.53) \quad \frac{\mathbf{R}}{R^n} - \frac{\mathbf{R}_h}{R_h^n} = \frac{\mathbf{R} - \mathbf{R}_h}{R^n} + \mathbf{R}_h \frac{(R_h - R)(R_h^{n-1} + R_h^{n-2}R + \dots + R^{n-1})}{R_h^n R^n}.$$

Using the bounds

$$(2.54) \quad |R - R_h| \leq |\mathbf{R} - \mathbf{R}_h| \leq ch^2$$

and $R_h \geq ch$, a consequence of (2.49), we conclude

$$(2.55) \quad \left| \frac{R}{R_h} - 1 \right| \leq ch,$$

which yields

$$(2.56) \quad \begin{cases} R_h \geq cR, \\ R \geq cR_h, \end{cases}$$

with a constant $c > 0$ not depending on h , L or K .

As a result, the estimate (2.50) follows in a straightforward way from (2.53). \square

Now we come to

Lemma 2.10. *There exists a constant c independent of h such that*

$$(2.57^1) \quad |e_h^{2,1}| \leq ch^{1/2} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}$$

for all μ_h in M_h and p_h in V_h .

Proof. Lemma 2.9 gives

$$(2.58) \quad |e_h^{2,1}| \leq ch \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{N}_K} \int_K \int_L \frac{|p(y_h) - p(x_h)|}{|y_h - x|^{n-1}} |\mu_K| dL(y_h) dK(x_h).$$

Repetition of the arguments used to prove (2.56) yields $|y_h - x| \geq c|y - x|$. Since p is constant along the normal \mathbf{n} to Γ (i.e., $p(y) = p(y_h)$ if $y = \mathcal{P}y_h$), there follows

$$(2.59) \quad \begin{aligned} |e_h^{2,1}| &\leq ch \int_{\Gamma} \int_{\Gamma} \frac{|p(y) - p(x)|}{|y - x|^{n-1}} |\mu(x)| d\Gamma(y) d\Gamma(x) \\ &= ch \int_{\Gamma} \int_{\Gamma} \frac{|p(y) - p(x)|}{|y - x|^{n-1-\alpha/2}} \frac{|\mu(x)|}{|y - x|^{\alpha/2}} d\Gamma(y) d\Gamma(x). \end{aligned}$$

If α is chosen so that $n - 3 < \alpha < n - 1$, the Cauchy-Schwarz inequality gives

$$(2.60) \quad \begin{aligned} |e_h^{2,1}| &\leq ch \left(\int_{\Gamma} \int_{\Gamma} \frac{|p(y) - p(x)|^2}{|y - x|^{2(n-1)-\alpha}} d\Gamma(y) d\Gamma(x) \right)^{1/2} \\ &\quad \cdot \left(\int_{\Gamma} \int_{\Gamma} \frac{|\mu(x)|^2}{|y - x|^{\alpha}} d\Gamma(y) d\Gamma(x) \right)^{1/2}. \end{aligned}$$

Note that

$$(2.61) \quad \left(\int_{\Gamma} \int_{\Gamma} \frac{|\mu(x)|^2}{|y - x|^{\alpha}} d\Gamma(y) d\Gamma(x) \right)^{1/2} \leq c_{\alpha} \|\mu\|_{0,\Gamma}$$

with

$$(2.62) \quad c_{\alpha} := \left(\sup_{x \in \Gamma} \int_{\Gamma} \frac{d\Gamma(y)}{|x - y|^{\alpha}} \right)^{1/2}.$$

Setting $n - 1 - \alpha = 2s$, if $0 < s < 1$ (i.e., $\alpha > n - 3$), one gets the bound (cf. Lions and Magenes [23])

$$(2.63) \quad \left(\int_{\Gamma} \int_{\Gamma} \frac{|p(y) - p(x)|^2}{|y - x|^{(n-1)+2s}} d\Gamma(y) d\Gamma(x) \right)^{1/2} \leq \|p\|_{s,\Gamma},$$

which yields

$$(2.64) \quad |e_h^{2,1}| \leq c_s h \|p\|_{s,\Gamma} \|\mu\|_{0,\Gamma} \quad \text{for all } s \text{ with } 0 < s < 1.$$

Using the inverse inequality, one can write

$$(2.65) \quad |e_h^{2,1}| \leq c_s h^{1/2} \|p\|_{s,\Gamma} \|\mu\|_{-1/2,\Gamma} \quad \text{for all } s \text{ with } 0 < s < 1.$$

It suffices to take $s = 1/2$, thus $\alpha = n - 2$, to complete the proof of the lemma. \square

Remark 2.11. In the bidimensional case (i.e., $n = 2$), we have $\alpha = 0$. This makes it possible to obtain directly the estimate (2.58) by means of the Cauchy-Schwarz inequality (cf. Bendali [2]). This is no longer possible for three-dimensional problems.

The next lemma gives an estimate for the term $e_h^{2,2}$.

Lemma 2.12. *There exists a constant c independent of h such that*

$$(2.57^2) \quad |e_h^{2,2}| \leq ch \|p_h\|_{1/2} \|\mu_h\|_{-1/2}$$

for all μ_h in M_h and p_h in V_h .

Proof. Since p_K is in \mathbb{P}_1 , the estimate of the term $e_h^{2,2}$ is obtained as follows:

$$(2.66) \quad \begin{aligned} e_h^{2,2} &= \sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \mathbf{grad}_K p_K \cdot \mathbf{R}_h \bar{\mu}_K dK(y_h) dK(x_h) \\ &= \sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \mathbf{grad}_K p_K \cdot \mathbf{R} \bar{\mu}_K dK(y_h) dK(x_h) \\ &\quad + \sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \mathbf{grad}_K p_K \cdot (\mathbf{R}_h - \mathbf{R}) \bar{\mu}_K dK(y_h) dK(x_h). \end{aligned}$$

Here, \mathbf{R} and \mathbf{R}_h are defined as in (2.51), and \mathbf{n}_K in the same way as \mathbf{n}_L .

Since $\mathbf{grad}_K p_K \cdot (\mathbf{R}_h - \mathbf{R})$ is constant as a function of y_h on K , there follows

$$(2.67) \quad \begin{aligned} &\sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \mathbf{grad}_K p_K \cdot (\mathbf{R}_h - \mathbf{R}) \bar{\mu}_K dK(y_h) dK(x_h) \\ &= \sum_{K \in \mathcal{T}_h} \int_K \bar{\mu}_K \mathbf{grad}_K p_K \cdot (\mathbf{R}_h - \mathbf{R}) \left(\int_K \partial_{\mathbf{n}_K} E(x, y_h) dK(y_h) \right) dK(x_h). \end{aligned}$$

The inner integral can be calculated by Gauss's theorem:

$$(2.68) \quad \int_K \partial_{\mathbf{n}_K} E(x, y_h) dK(y_h) = 2 \frac{\omega(K, x)}{\omega_n},$$

where $\omega(K, x)$ is the angle (resp. the solid angle) under which the element K is viewed from the point x for $n = 2$ (resp. $n = 3$).

Using (2.54), we can bound the term (2.67) by

$$(2.69) \quad ch^2 \sum_{K \in \mathcal{T}_h} \int_K |\mathbf{grad}_K p_K| |\mu_K| dK(x_h).$$

The Cauchy-Schwarz inequality and the inverse property of Lemma 1.5 yield

$$(2.70) \quad \left| \sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \operatorname{grad}_K p_K \cdot (\mathbf{R}_h - \mathbf{R}) \bar{\mu}_K dK(y_h) dK(x_h) \right| \leq ch \|p_h\|_{1/2} \|\mu_h\|_{-1/2}.$$

The last term to estimate in $e_h^{2,2}$ becomes

$$(2.71) \quad \tilde{e}_h^{2,2} := \sum_{K \in \mathcal{T}_h} \int_K \int_K \partial_{\mathbf{n}_K} E(x, y_h) \operatorname{grad}_K p_K \cdot \mathbf{R} \bar{\mu}_K dK(y_h) dK(x_h).$$

A straightforward calculation gives

$$(2.72) \quad \begin{aligned} \partial_{\mathbf{n}_K} E(x, y_h) \operatorname{grad}_K p_K \cdot \mathbf{R} &= -\frac{2}{\omega_n} \frac{\mathbf{R} \cdot \mathbf{n}_K}{R_h^n} \operatorname{grad}_K p_K \cdot \mathbf{R} \\ &= \mathbf{R} \cdot \mathbf{n}_K \operatorname{grad}_K p_K \cdot \operatorname{grad}_K E(x, y_h). \end{aligned}$$

Since $\mathbf{R} \cdot \mathbf{n}_K$ is constant when x is fixed and y_h runs over K , the term $\tilde{e}_h^{2,2}$ may be written as

$$(2.73) \quad \tilde{e}_h^{2,2} = \sum_{K \in \mathcal{T}_h} \int_K \bar{\mu}_K \mathbf{R} \cdot \mathbf{n}_K \operatorname{grad}_K p_K \cdot \left(\int_K \operatorname{grad}_K E(x, y_h) dK(y_h) \right) dK(x_h).$$

Two cases need to be distinguished.

(i) Two-dimensional case ($n = 2$). In this case, K is a segment having A_K and B_K as endpoints. If $\boldsymbol{\tau}_K$ denotes the unit tangent vector oriented from A_K toward B_K , an elementary theorem of integral calculus gives

$$(2.74) \quad \begin{aligned} \int_K \operatorname{grad}_K E(x, y_h) dK(y_h) &= (E(x, B_K) - E(x, A_K)) \boldsymbol{\tau}_K \\ &= -\frac{1}{\pi} (\ln(|x - B_K|) - \ln(|x - A_K|)) \boldsymbol{\tau}_K. \end{aligned}$$

Since $|\mathbf{R} \cdot \mathbf{n}_K|$ is the distance of the point x to the segment K , we have the bound

$$(2.75) \quad |\tilde{e}_h^{2,2}| \leq \frac{2}{\pi} \sum_{K \in \mathcal{T}_h} \int_K |\mu_K| |\operatorname{grad}_K p_K| |\mathbf{R} \cdot \mathbf{n}_K \ln(|\mathbf{R} \cdot \mathbf{n}_K|)| dK(x_h).$$

Since $|\mathbf{R} \cdot \mathbf{n}_K| \leq ch^2$ and $|\mathbf{R} \cdot \mathbf{n}_K|^\varepsilon \ln(|\mathbf{R} \cdot \mathbf{n}_K|)$ is bounded for any fixed $\varepsilon > 0$, it follows that for all ε such that $0 < \varepsilon < 1/2$ there exists $h_\varepsilon^* > 0$ and a constant c_ε not depending on $0 < h < h_\varepsilon^*$ such that

$$(2.76) \quad |\tilde{e}_h^{2,2}| \leq c_\varepsilon h^{1-\varepsilon} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}.$$

(ii) Three-dimensional case ($n = 3$). Following the method initiated by De Hoop (cf. Van Herk [36]), we can express

$$(2.77) \quad \int_K \operatorname{grad}_{y_h} \frac{1}{|x - y_h|} dK(y_h) = \mathbf{n}_K \omega(K, x) + \sum_{m=1}^3 \nu_m^K \ln \frac{R_{m+1}^K + \mathbf{R}_{m+1}^K \cdot \boldsymbol{\tau}_m^K}{R_m^K + \mathbf{R}_m^K \cdot \boldsymbol{\tau}_m^K},$$

where A_1, A_2 , and A_3 are the three vertices of the triangle K , $\mathbf{R}_m^K := A_m - x$, $R_m^K := |\mathbf{R}_m^K|$, $\boldsymbol{\tau}_m^K := (A_{m+1} - A_m)/|A_{m+1} - A_m|$, for $m = 1, 2, 3$; see Figure 2. The index convention that $m + 1 = 1$ when $m = 3$ is adopted.

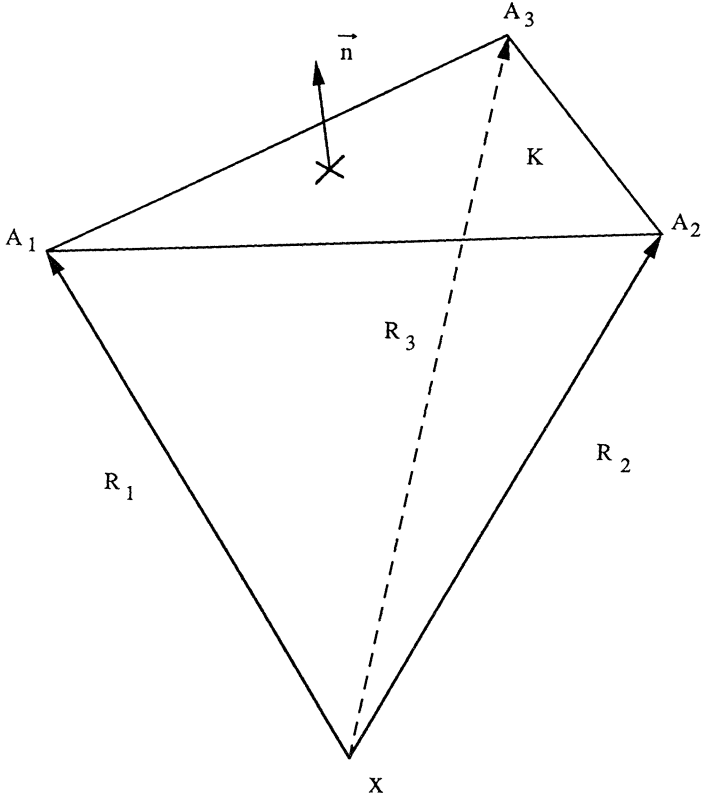


FIGURE 2

An elementary calculation gives

$$(2.78) \quad \ln \frac{R_{m+1}^K + \mathbf{R}_{m+1}^K \cdot \boldsymbol{\tau}_m^K}{R_m^K + \mathbf{R}_m^K \cdot \boldsymbol{\tau}_m^K} = -\ln(R_{m+1}^K - \mathbf{R}_{m+1}^K \cdot \boldsymbol{\tau}_m^K) + \ln(R_m^K - \mathbf{R}_m^K \cdot \boldsymbol{\tau}_m^K),$$

thus yielding

$$(2.79) \quad \forall \varepsilon > 0, \exists c_\varepsilon : |\mathbf{n}_K \cdot \mathbf{R}|^\varepsilon \left| \ln \frac{R_{m+1}^K + \mathbf{R}_{m+1}^K \cdot \boldsymbol{\tau}_m^K}{R_m^K + \mathbf{R}_m^K \cdot \boldsymbol{\tau}_m^K} \right| \leq c_\varepsilon.$$

This completes the proof of the lemma. \square

It remains to estimate the term $e_h^{2,3}$. For this, we need the following lemma.

Lemma 2.13. *For all K in \mathcal{T}_h and L in \mathcal{A}_K , and for all x_h in the interior of K , we have*

$$(2.80) \quad \begin{cases} \omega(L, x) \leq ch, \\ \omega(L, x_h) \leq ch. \end{cases}$$

The notation is that introduced for K in (2.68).

Proof. This is a consequence of well-known facts about error estimates for the derivatives in finite element interpolation. The details of how to obtain these estimates are not given. However, note that these estimates are obvious in

the two-dimensional case, and a consequence of the classical formula explicitly giving the expression of $\omega(L, x)$ and $\omega(L, x_h)$ in the three-dimensional case (cf., e.g., Van Herk [36]).

The last estimate is given by the following lemma.

Lemma 2.14. *There exists a constant c independent of h such that*

$$(2.81) \quad |e_h^{2,3}| \leq ch^{1/2} \|p_h\|_{1/2} \|\mu_h\|_{-1/2}$$

for all μ_h in M_h and p_h in V_h .

Proof. Consider two elements in \mathcal{T}_h such that L and K are different, but have at least one common point M . Recall the notations $p_K := p_h|_K$ and $p_L := p_h|_L$. Arguments previously used give

$$(2.82) \quad \begin{aligned} p_L(y_h) - p_K(x_h) &= p_L(y_h) - p_L(M) + p_K(M) - p_K(x_h) \\ &= (y_h - M) \cdot \text{grad}_L p_L + (M - x_h) \cdot \text{grad}_K p_K \\ &= \mathbf{R} \cdot \text{grad}_L p_L + (x - M) \cdot \text{grad}_L p_L + (M - x_h) \cdot \text{grad}_K p_K. \end{aligned}$$

Observing that $(x - M) \cdot \text{grad}_L p_L + (M - x_h) \cdot \text{grad}_K p_K$ is constant in y_h , we get

$$(2.83) \quad \begin{aligned} e_h^{2,3} &= \frac{2}{\omega_n} \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{A}_K} \int_K ((x - M) \cdot \text{grad}_L p_L + (M - x_h) \cdot \text{grad}_K p_K) \\ &\quad \cdot (\omega(L, x) - \omega(L, x_h)) \bar{\mu}_K dK(x_h) \\ &\quad - \frac{2}{\omega_n} \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{A}_K} \int_K \int_L \left(\frac{\mathbf{R}}{R^n} - \frac{\mathbf{R}_h}{R_h^n} \right) \cdot \mathbf{n}_L \mathbf{R} \cdot \text{grad}_L p_L \bar{\mu}_K dL(y_h) dK(x_h). \end{aligned}$$

Clearly,

$$(2.84) \quad \begin{cases} |x - M| \leq ch, \\ |x_h - M| \leq ch. \end{cases}$$

It follows from Lemma 2.13 that the first term of the right-hand side of (2.83) can be bounded by

$$(2.85) \quad ch^2 \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{A}_K} \int_K (|\text{grad}_K p_K| + |\text{grad}_L p_L|) |\mu_K| dL(y_h) dK(x_h).$$

The uniform regularity of the triangulation \mathcal{T}_h (cf. (1.29)) insures that there exists a positive integer N not depending on K or h that bounds the number of elements of \mathcal{A}_K , i.e.,

$$(2.86) \quad \text{Card} \mathcal{A}_K \leq N.$$

Clearly, arguments used above enable us to bound the expression (2.85) by terms of the form $ch \|\mu_h\|_{-1/2} \|p_h\|_{1/2}$.

Introducing the common point M in the second term of (2.83), thus writing

$$(2.87) \quad \frac{\mathbf{R} \cdot \mathbf{n}_L}{R^n} \mathbf{R} \cdot \text{grad}_L p_L = (M - x) \cdot \mathbf{n}_L \frac{\mathbf{R}}{R^n} \cdot \text{grad}_L p_L,$$

$$\begin{aligned}
 (2.88) \quad & \frac{\mathbf{R}_h \cdot \mathbf{n}_L}{R_h^n} \mathbf{R} \cdot \text{grad}_L p_L \\
 & = (M - x_h) \cdot \mathbf{n}_L \frac{\mathbf{R}_h}{R_h^n} \cdot \text{grad}_L p_L + \frac{\mathbf{R}_h \cdot \mathbf{n}_L}{R_h^n} (\mathbf{R} - \mathbf{R}_h) \cdot \text{grad}_L p_L
 \end{aligned}$$

helps estimating the term

$$\begin{aligned}
 (2.89) \quad \tilde{\epsilon}_h^{2,3} := & \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{O}_K} \left\{ \int_K (M - x) \cdot \mathbf{n}_L \bar{\mu}_K \text{grad}_L p_L \cdot \mathbf{g}_L(x) dK(x_h) \right. \\
 & \left. - \int_K (M - x_h) \cdot \mathbf{n}_L \bar{\mu}_K \text{grad}_L p_L \cdot \mathbf{g}_L(x_h) dK(x_h) \right\} \\
 & - \frac{2}{\omega_n} \sum_{K \in \mathcal{T}_h} \sum_{L \in \mathcal{O}_K} \int_K (\mathbf{R} - \mathbf{R}_h) \cdot \text{grad}_L p_L \bar{\mu}_K \omega(L, x_h) dK(x_h).
 \end{aligned}$$

The vector-valued function defined for z in \mathbb{R}^n by

$$\mathbf{g}_L(z) := \int_L \text{grad}_L E(z, y_h) dL(y_h)$$

is introduced to make the preceding formulae more concise. Using the estimates

$$(2.90) \quad \begin{cases} |(M - x) \cdot \mathbf{n}_L| \leq ch^2, \\ |(M - x_h) \cdot \mathbf{n}_L| \leq ch^2, \\ |\mathbf{R} - \mathbf{R}_h| \leq ch^2, \end{cases}$$

the De Hoop method (cf. (2.77)) and previously invoked arguments, we get

$$(2.91) \quad |\tilde{\epsilon}_h^{2,3}| \leq c_\varepsilon h^{1-\varepsilon} \|\mu_h\|_{-1/2} \|p_h\|_{1/2}.$$

Combining all the estimates given in the previous lemmas will prove Theorem 2.5, and hence the main results given in Theorem 2.1.

Remark 2.15. The convergence rate $h^{1/2}$ of the discrete solution to the exact one seems to be somewhat weak. This is due to the fact that some estimates are given in the norm of $H^{-1/2}(\Gamma)$, because this space is necessary for the coerciveness of the sesquilinear forms. However, in similar situations (cf. Bendali [2]), numerical results show that the convergence rate in the stronger norm of $L^2(\Gamma)$ is of the order h . Nevertheless, the $h^{1/2}$ -estimate may be considered as a consistency result, which makes the numerical analysis of the exact problem possible: it establishes existence, uniqueness, stability of the solution of the discrete problem, and its convergence to the solution of the continuous problem. The technical assumptions on the triangulation (cf. (1.28), (1.29), and (1.30)) seem to be needed only for purposes of theoretical asymptotic analysis and are not essential in practical computations.

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